

# Quantum mechanics of photons and Maxwell's equations

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## Motto

“The main principle of the present work is the idea that, since matter and light both possess the dual characters of particle and wave, a similar mathematical treatment should be applied to both, and that this has not been yet done as fully as should be possible”.

Charles G. Darwin 1932

# The Schrödinger equation for photons

Weyl equation for (massless) neutrinos (spin 1/2)

$$i\hbar\partial_t\psi = c(\boldsymbol{\sigma}\cdot\mathbf{p})\psi \quad \text{or} \quad i\partial_t\psi = -ic(\boldsymbol{\sigma}\cdot\nabla)\psi$$

Analogous equation for photons (spin 1)

$$i\partial_t\mathbf{F} = -ic(\mathbf{S}\cdot\nabla)\mathbf{F}$$

$$\mathbf{S}\cdot\nabla = \begin{bmatrix} 0 & -i\partial_z & i\partial_y \\ i\partial_z & 0 & -i\partial_x \\ -i\partial_y & i\partial_x & 0 \end{bmatrix} = i\nabla\times$$

## Splitting into real and imaginary parts

$$i\partial_t \mathbf{F} = c \nabla \times \mathbf{F}$$

$$\mathbf{F} = \Re(\mathbf{F}) + i\Im(\mathbf{F})$$

$$\partial_t \Re(\mathbf{F}) = c \nabla \times \Im(\mathbf{F}) \quad \partial_t \Im(\mathbf{F}) = -c \nabla \times \Re(\mathbf{F})$$

$$\Re(\mathbf{F}) = \sqrt{\epsilon} \mathbf{E} = \frac{\mathbf{D}}{\sqrt{\epsilon}} \quad \Im(\mathbf{F}) = \sqrt{\mu} \mathbf{H} = \frac{\mathbf{B}}{\sqrt{\mu}}$$

$$\partial_t \mathbf{D} = \nabla \times \mathbf{H} \quad \partial_t \mathbf{B} = -\nabla \times \mathbf{E}$$

# Relativistic quantum mechanics of photons

General solution of the Schrödinger equation  
for photons

$$\mathbf{F}(\mathbf{r}, t) = \sqrt{\hbar c} \int \frac{d^3k}{(2\pi)^{3/2}} \mathbf{e}(\mathbf{k}) \left[ f_L(\mathbf{k}) e^{-i\omega_{\mathbf{k}}t + i\mathbf{k}\cdot\mathbf{r}} + f_R^*(\mathbf{k}) e^{i\omega_{\mathbf{k}}t - i\mathbf{k}\cdot\mathbf{r}} \right]$$

Polarization vector obeys the (Maxwell) equation

$$-i\omega \mathbf{e}(\mathbf{k}) = c\mathbf{k} \times \mathbf{e}(\mathbf{k}) \quad \mathbf{e}^* \cdot \mathbf{e} = 1$$

The amplitudes  $f_L(\mathbf{k})$  and  $f_R(\mathbf{k})$  are  
the wave functions in momentum space

## Quantum operators in $k$ -space

Generators of Poincaré transformations

The transformations  $F'_i(\mathbf{r}', t') = O_i^j F_j(\mathbf{r}, t)$  must preserve the form of the photon wave equation

Time translation: Energy =  $\hbar\omega$

Space translation: Momentum =  $\hbar\mathbf{k}$

Rotation: Angular momentum =  $i\hbar\mathbf{k} \times \mathcal{D}_{\mathbf{k}} + \hat{\chi}\hbar\mathbf{k}/k$

Lorentz transformation: Boost =  $i\hbar\omega\mathcal{D}_{\mathbf{k}}$

Helicity operator  $\hat{\chi}$  takes on two values  $\pm 1$

$$\mathcal{D}_{\mathbf{k}} = \nabla_{\mathbf{k}} - i\hat{\chi}\boldsymbol{\alpha}(\mathbf{k}) \quad \nabla_{\mathbf{k}} \times \boldsymbol{\alpha}(\mathbf{k}) = -\mathbf{k}/k^3$$

## Quantum-classical correspondence

QM average values agree with classical expressions

$$\text{Energy} = \langle \hbar\omega \rangle = \int d^3r [\mathbf{D}^2/2\epsilon + \mathbf{B}^2/2\mu]$$

$$\text{Momentum} = \langle \hbar\mathbf{k} \rangle = \int d^3r [\mathbf{D} \times \mathbf{B}]$$

$$\begin{aligned} \text{Angular momentum} &= \langle i\hbar\mathbf{k} \times \mathcal{D}_{\mathbf{k}} + \hat{\chi}\hbar\mathbf{k}/k \rangle \\ &= \int d^3r [\mathbf{r} \times (\mathbf{D} \times \mathbf{B})] \end{aligned}$$

$$\begin{aligned} \text{Lorentz transformation: Boost} &= \langle i\hbar\omega\mathcal{D}_{\mathbf{k}} \rangle \\ &= \int d^3r \mathbf{r} [\mathbf{D}^2/2\epsilon + \mathbf{B}^2/2\mu] \end{aligned}$$

## Second quantization

Quantized electromagnetic field operator

$$\hat{\mathbf{F}}(\mathbf{r}, t) = \sqrt{\hbar c} \int \frac{d^3 k}{(2\pi)^{3/2}} \mathbf{e}(\mathbf{k}) \left[ a_L(\mathbf{k}) e^{-i\omega_{\mathbf{k}} t + i\mathbf{k} \cdot \mathbf{r}} + a_R^\dagger(\mathbf{k}) e^{i\omega_{\mathbf{k}} t - i\mathbf{k} \cdot \mathbf{r}} \right]$$

Photons **do not have** a conserved quantum number  
(charge, lepton number, etc.)

Formally, right-handed and left-handed photons  
are in the particle-antiparticle relation but  
we can make **all their superpositions** that create  
photon states with arbitrary polarization

$$|\Psi_{\text{one photon}}\rangle = \int \frac{d^3 k}{k} \left[ f_L(\mathbf{k}) a_L^\dagger(\mathbf{k}) + f_R(\mathbf{k}) a_R^\dagger(\mathbf{k}) \right] |0\rangle$$



How come classical EM field?

How come Maxwell's equations?

Key property: Number of photons  $N$

$$N = \frac{\text{Power} \times \text{Time}}{\text{Photon energy}} = 7.5 \times 10^{31} \frac{P[\text{in Watt}] \times T[\text{in Sec}]}{\nu[\text{in Hertz}]}$$

Small WiFi router (50mW) operating at 2.4GHz  
sends  $3 \times 10^{22}$  photons per second

## Coherent states

Note that the average field in any state with fixed number of photons vanishes  $\langle \Psi_N | \hat{F} | \Psi_N \rangle = 0$

It is obvious that one cannot precisely control  $N$  at the level of  $10^{22}$

Randomly produced photons are characterized by the Poisson distribution  $\langle N \rangle^k / k! e^{-\langle N \rangle}$

The Poissonian quantum-mechanical state is:

$$|\alpha\rangle = \sum_1^{\infty} \frac{\alpha^k}{\sqrt{k!}} |k\rangle \quad |\alpha|^2 = \langle N \rangle$$

This state is called **coherent state**

## Average field

Assume that a device (say a router) produces photons characterized by the creation operator

$$a_f^\dagger = \int \frac{d^3k}{k} \left[ f_L(\mathbf{k}) a_L^\dagger(\mathbf{k}) + f_R(\mathbf{k}) a_R^\dagger(\mathbf{k}) \right]$$

The classical electromagnetic field is the average value obtained from the complex average value  $\langle \hat{\mathbf{F}} \rangle_f$  of  $\hat{\mathbf{F}}$  calculated in the coherent state corresponding to  $a_f^\dagger$

$$\langle \hat{\mathbf{F}} \rangle_f = \sqrt{\langle N \rangle \hbar c} \int \frac{d^3k}{(2\pi)^{3/2}} \mathbf{e}(\mathbf{k}) \left[ f_L(\mathbf{k}) e^{-i\omega_{\mathbf{k}}t + i\mathbf{k}\cdot\mathbf{r}} + f_R^*(\mathbf{k}) e^{i\omega_{\mathbf{k}}t - i\mathbf{k}\cdot\mathbf{r}} \right]$$

## Closing the argument

In classical electrodynamics classical sources produce **classical electromagnetic field**

What state  $|\Psi\rangle$  of the **quantum electromagnetic field** is produced by a classical current  $\mathbf{J}^\mu(\mathbf{r}, t)$ ?

The answer is obtained from the formula

$$|\Psi\rangle = \mathbf{T} \exp \left( -i \int d^4x \hat{\mathbf{A}}_\mu(\mathbf{r}, t) \mathbf{J}^\mu(\mathbf{r}, t) \right) |0\rangle$$

The state  $|\Psi\rangle$  is a coherent state and the average field in this state is **the same** as the one **obtained from the classical theory!**

## Summary

Maxwell's equations can be derived from quantum mechanics of photons in the classical limit

The classical limit means here not  $\hbar \rightarrow 0$  but

a very large average number of photons

obeying the Poisson distribution

Classical fields are identified as

expectation values of the quantum field operators